

Feynman Rules - The Standard Model

The Lagrangian density for the Standard Model, after spontaneous symmetry breaking in the *unitary gauge* for the EW part and ignoring the gauge fixing for QCD, is given by

$$\begin{aligned}
 \mathcal{L}_{\text{sm}} = & \frac{i}{2} \sum_A \bar{u}_A \not{\partial} u_A + \frac{i}{2} \sum_A \bar{d}_A \not{\partial} d_A + \frac{i}{2} \sum_A \bar{\ell}_A \not{\partial} \ell_A + \frac{i}{2} \sum_A \bar{\nu}_A \not{\partial} \nu_A + \text{h.c.} + \frac{1}{2} \partial_\mu h \partial^\mu h \\
 & - \frac{1}{4} (\partial_\mu Z_\nu - \partial_\nu Z_\mu - ig \cos \theta_W (W_\mu^+ W_\nu^- - W_\mu^- W_\nu^+))^2 \\
 & - \frac{1}{4} (\partial_\mu A_\nu - \partial_\nu A_\mu - ie (W_\mu^+ W_\nu^- - W_\mu^- W_\nu^+))^2 \\
 & - \frac{1}{2} |\partial_\mu W_\nu^+ - \partial_\nu W_\mu^+ + ig \cos \theta_W (W_\mu^+ Z_\nu - W_\nu^+ Z_\mu) + ie (W_\mu^+ A_\nu - W_\nu^+ A_\mu)|^2 \\
 & - \sum_A m_{u_A} \bar{u}_A u_A - \sum_A m_{d_A} \bar{d}_A d_A - \sum_A m_{\ell_A} \bar{\ell}_A \ell_A - \frac{1}{2} m_h^2 h^2 + m_W^2 W_\mu^+ W^{-\mu} + \frac{1}{2} m_Z^2 Z_\mu Z^\mu \\
 & - \frac{g}{\sqrt{2}} \sum_A \bar{\nu}_A \gamma^\mu P_L \ell_A W_\mu^+ + \text{h.c.} - \frac{g}{\sqrt{2}} \sum_{A,B} \bar{d}_A \gamma^\mu P_L V_{AB} u_B W_\mu^+ + \text{h.c.} \\
 & - \sum_f Q_f e \bar{f} \gamma^\mu f A_\mu - \frac{g}{\cos \theta_W} \sum_f \bar{f} \gamma^\mu (v_f - a_f \gamma_5) f Z_\mu \\
 & - \frac{1}{2} g \sum_A \frac{m_{u_A}}{m_W} h \bar{u}_A u_A - \frac{1}{2} g \sum_A \frac{m_{d_A}}{m_W} h \bar{d}_A d_A - \frac{1}{2} g \sum_A \frac{m_{\ell_A}}{m_W} h \bar{\ell}_A \ell_A - \frac{3}{2} g \frac{m_h^2}{m_W} h^3 - \frac{3}{4} g^2 \frac{m_h^2}{m_W^2} h^4 \\
 & + \frac{1}{4} g^2 h^2 W_\mu^+ W^{-\mu} + \frac{1}{8} \frac{g^2}{\cos^2 \theta_W} h^2 Z_\mu Z^\mu + m_W g h W_\mu^+ W^{-\mu} + \frac{1}{2} \frac{g}{\cos \theta_W} m_W h Z_\mu Z^\mu \\
 & - \frac{1}{4} G_{\mu\nu}^a G^{\mu\nu a}
 \end{aligned} \tag{1}$$

where A is the generation index $A = \{1, 2, 3\}$, ℓ_A are the leptons $\ell_A = \{e, \mu, \tau\}$, ν_A are the neutrinos $\nu_A = \{\nu_e, \nu_\mu, \nu_\tau\}$, u_A are the up-type quarks $u_A = \{u, c, t\}$, d_A are the down-type quarks $d_A = \{d, s, b\}$, f is the entire set of fermions.

Feynman Rules

Here we give the Feynman rules for the scattering amplitude \mathcal{M} ,

$$i\mathcal{M} = \text{sum of all connected, amputated diagrams,}$$

where the diagrams are evaluated according to the following rules:

- Draw all topologically distinct diagrams at a given order;
- For each internal gauge boson line, attach a propagator

$$\begin{array}{c} \gamma \\ \mu \text{ --- } \text{---} \text{---} \text{---} \nu \\ \xrightarrow{p} \end{array} = \frac{-ig_{\mu\nu}}{p^2 + i\epsilon};$$

$$\mu \begin{array}{c} \text{Z} \\ \text{~~~~~} \\ \xrightarrow{p} \end{array} \nu = \frac{-i}{p^2 - m_Z^2 + i\epsilon} \left(g_{\mu\nu} - \frac{p_\mu p_\nu}{m_Z^2} \right);$$

$$\mu \begin{array}{c} \text{W} \\ \text{~~~~~} \\ \xrightarrow{p} \end{array} \nu = \frac{-i}{p^2 - m_W^2 + i\epsilon} \left(g_{\mu\nu} - \frac{p_\mu p_\nu}{m_W^2} \right);$$

- For each internal fermion line, attach a propagator

$$\begin{array}{c} \longrightarrow \\ \xrightarrow{p} \end{array} = \frac{i(\not{p} + m_f)}{p^2 - m_f^2 + i\epsilon};$$

- For each internal higgs boson line, attach a propagator

$$\begin{array}{c} \text{-----} \\ \xrightarrow{p} \end{array} = \frac{i}{p^2 - m_h^2 + i\epsilon};$$

- For each fermion-neutral gauge boson vertex, assign

$$\begin{array}{c} \psi_f \\ \nearrow \\ \mu \text{ } \gamma \\ \text{~~~~~} \\ \bullet \\ \searrow \\ \psi_f \end{array} = -iQ_f e \gamma^\mu;$$

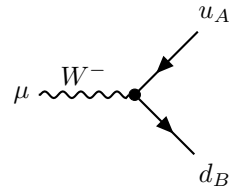
$$\begin{array}{c} \psi_f \\ \nearrow \\ \mu \text{ } \text{Z} \\ \text{~~~~~} \\ \bullet \\ \searrow \\ \psi_f \end{array} = -i \frac{g}{\cos \theta_W} \gamma^\mu (v_f - a_f \gamma_5);$$

$$v_f = \frac{1}{2} T_f^3 - Q_f \sin^2 \theta_W, \quad a_f = \frac{1}{2} T_f^3.$$

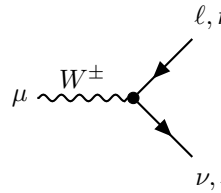
Note: $T_\ell^3 = -1/2$, $T_\nu^3 = +1/2$, $Q_\ell = -1$, $Q_\nu = 0$.

- For each fermion-charged gauge boson vertex, assign

$$\begin{array}{c} d_B \\ \nearrow \\ \mu \text{ } \text{W}^+ \\ \text{~~~~~} \\ \bullet \\ \searrow \\ u_A \end{array} = -i \frac{g}{\sqrt{2}} \gamma^\mu \frac{1}{2} (1 - \gamma_5) V_{AB};$$

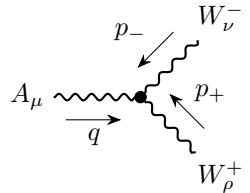


$$= -i \frac{g}{\sqrt{2}} \gamma_\mu \frac{1}{2} (1 - \gamma_5) V_{AB}^* ;$$

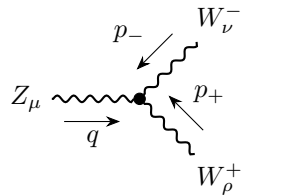


$$= -i \frac{g}{\sqrt{2}} \gamma_\mu \frac{1}{2} (1 - \gamma_5) ;$$

- For each triple gauge vertex, assign

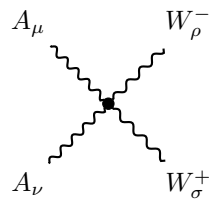


$$= ie \left[(q - p_-)^\rho g^{\mu\nu} + (p_- - p_+)^\mu g^{\nu\rho} + (p_+ - q)^\nu g^{\mu\rho} \right] ;$$

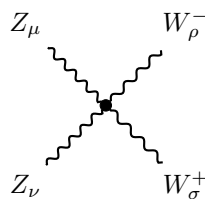


$$= -ig \cos \theta_W \left[(q - p_-)^\rho g^{\mu\nu} + (p_- - p_+)^\mu g^{\nu\rho} + (p_+ - q)^\nu g^{\mu\rho} \right] ;$$

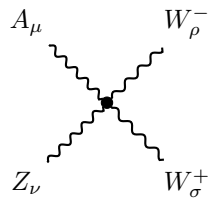
- For each quartic gauge vertex, assign



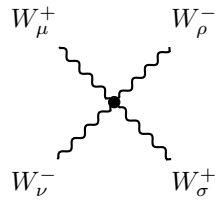
$$= -ie^2 \left[2g_{\mu\nu}g_{\rho\sigma} - g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma} \right] ;$$



$$= -ig^2 \cos^2 \theta_W \left[2g_{\mu\nu}g_{\rho\sigma} - g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma} \right] ;$$

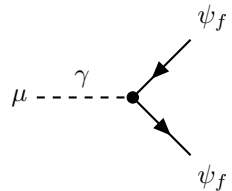


$$= -ieg \cos \theta_W \left[2g_{\mu\nu}g_{\rho\sigma} - g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma} \right];$$



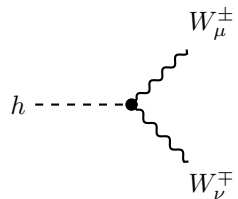
$$= ig^2 \left[2g_{\mu\nu}g_{\rho\sigma} - g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma} \right];$$

- For each fermion-Higgs vertex, assign

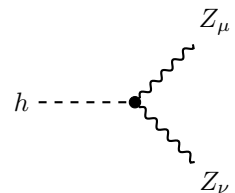


$$= -i \frac{g}{2} \frac{m_f}{m_W};$$

- For each gauge-Higgs three-vertex, assign

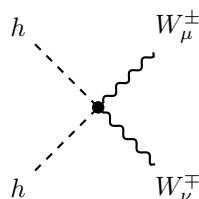


$$= igm_W g^{\mu\nu};$$

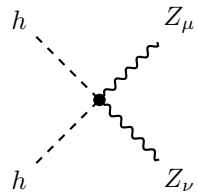


$$= i \frac{g}{\cos \theta_W} m_Z g^{\mu\nu};$$

- For each gauge-Higgs four-vertex, assign

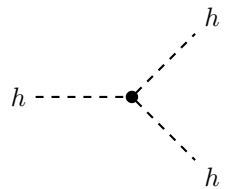


$$= -i \frac{g^2}{2} g^{\mu\nu};$$

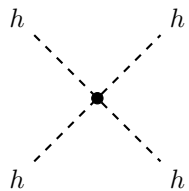


$$= -i \frac{1}{2} \frac{g^2}{\cos^2 \theta_W} g^{\mu\nu};$$

- For each Higgs self-interaction, assign

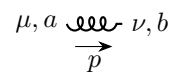


$$= -i \frac{3}{2} g \frac{m_h^2}{m_W};$$



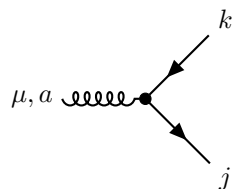
$$= -i \frac{3}{4} g^2 \frac{m_h^2}{m_W^2};$$

- For each internal gluon line, attach a propagator



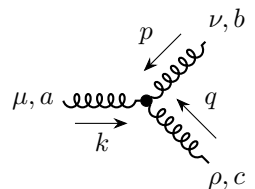
$$= \frac{-i \delta_{ab}}{p^2 + i\epsilon} \left(g_{\mu\nu} - (1 - \xi_G) \frac{p_\mu p_\nu}{p^2} \right);$$

- For each quark-gluon vertex, assign



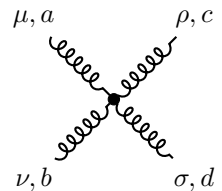
$$= -i g_s \gamma^\mu (T_a)_{jk};$$

- For each three-gluon vertex, assign



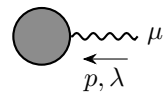
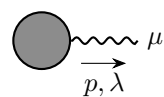
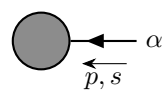
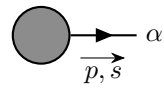
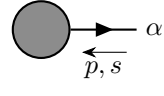
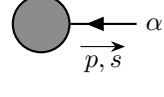
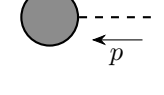
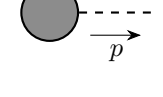
$$= -i g_s (-i f_{abc}) \left[(k - p)^\rho g^{\mu\nu} + (p - q)^\mu g^{\nu\rho} + (q - k)^\nu g^{\mu\rho} \right];$$

- For each four-gluon vertex, assign



$$= -ig_s^2 \left[\begin{aligned} & f_{abe} f_{cde} (g_{\mu\rho} g_{\nu\sigma} - g_{\mu\sigma} g_{\nu\rho}) \\ & + f_{ace} f_{bde} (g_{\mu\nu} g_{\rho\sigma} - g_{\mu\sigma} g_{\nu\rho}) \\ & + f_{ade} f_{bce} (g_{\mu\nu} g_{\rho\sigma} - g_{\mu\rho} g_{\nu\sigma}) \end{aligned} \right];$$

- For each external line, place the particle on the “mass-shell”, $p^2 = m_q^2$ for the quark and $p^2 = 0$ for the gluon, and attach a wavefunction factor

“incoming boson”		$= \epsilon_\mu(p, \lambda);$
“outgoing boson”		$= \epsilon_\mu^*(p, \lambda);$
“incoming lepton”		$= u_\alpha(p, s);$
“outgoing lepton”		$= \bar{u}_\alpha(p, s);$
“incoming anti-lepton”		$= \bar{v}_\alpha(p, s);$
“outgoing anti-lepton”		$= v_\alpha(p, s);$
“incoming Higgs”		$= 1;$
“outgoing Higgs”		$= 1;$

- Impose momentum conservation at each vertex;
- For each internal loop momentum k not fixed by momentum conservation, integrate $\int \frac{d^4k}{(2\pi)^4};$
- For each fermion loop, multiply the diagram by $(-1);$

- For each set of diagram which are only distinguished by interchanging two external fermion lines, multiply one of the diagrams by (-1) ;
- Multiply the contribution for each diagram by an appropriate symmetry factor \mathcal{S}^{-1} for identical particles.