

Radiative Corrections

We are now (finally) ready to compute the first-order radiative correction to the magnetic moment of the electron

$$a_e^{(1)} = F_2^{(1)}(0)$$

$$= \frac{1}{12m^2} \text{tr} \left[\left(m^2 \gamma_\mu - \not{P} \not{P} - \frac{3}{2} m \not{P}_\mu \right) V_1^\mu + \frac{m}{4} \left(\frac{\not{P}}{2} + m \right) [\gamma_\mu, \gamma_\nu] \left(\frac{\not{P}}{2} + m \right) \delta V_1^{\mu\nu} \right]$$

where $V_1^\mu(P) = \Lambda_1^\mu(P, 0)$

$$\delta V_1^{\mu\nu}(P) = \left. \frac{\partial \Lambda_1^\mu(P, \xi)}{\partial \xi_\nu} \right|_{\xi=0}$$

where $\Lambda_1^\mu(P, \xi) \equiv \Lambda_1^\mu(p', p)$ is the correction to the vertex function

$$\Lambda_1^\mu(p', p) = -e^2 \int \frac{d^4 k}{(2\pi)^4} \frac{-i^3 N^\mu(p', p, k)}{k^2 [(p'-k)^2 - m^2] [(p-k)^2 - m^2]}$$

with $N^\mu = \gamma_\nu [(p'-k) + m] \gamma^\mu [(p-k) + m] \gamma^\nu$

Let us now replace the kinematics

$$p' = \frac{1}{2}(P + \xi), \quad p = \frac{1}{2}(P - \xi)$$

S_2

$$\Lambda_1^{\wedge}(P, q) = -i e^2 \int \frac{d^4 k}{(2\pi)^4} \frac{N^{\wedge}(P, q, k)}{k^2 [((P+q)/2 - k)^2 - m^2] [(P-q)/2 - k)^2 - m^2]}$$

with

$$N^{\wedge}(P, q, k) = \gamma_{\alpha} [(P+q)/2 - k + m] \gamma^{\wedge} [(P-q)/2 - k + m] \gamma^{\alpha}$$

At $q=0$,

$$\Lambda_1^{\wedge}(P, 0) = -i e^2 \int \frac{d^4 k}{(2\pi)^4} \frac{N^{\wedge}(P, 0, k)}{k^2 [(P/2 - k)^2 - m^2]^2}$$

with $N^{\wedge}(P, 0, k) = \gamma_{\alpha} [P/2 - k + m] \gamma^{\wedge} [P/2 - k + m] \gamma^{\alpha}$

the derivative is given by

$$\frac{\partial \Lambda_1^{\wedge}}{\partial q_{\nu}} \Big|_{q=0} = -i e^2 \int \frac{d^4 k}{(2\pi)^4} \left\{ \frac{1}{k^2 [(P/2 - k)^2 - m^2]^2} \frac{\partial N^{\wedge}}{\partial q_{\nu}} \Big|_{q=0} + \frac{N^{\wedge}(P, 0, k)}{k^2} \times \frac{\partial}{\partial q_{\nu}} \left(\frac{1}{[(P+q)/2 - k)^2 - m^2] [(P-q)/2 - k)^2 - m^2]} \right) \Big|_{q=0} \right\}$$

Now, $\frac{\partial N^{\wedge}}{\partial q_{\nu}} \Big|_{q=0} = \gamma_{\alpha} \frac{\partial}{\partial q_{\nu}} [(P+q)/2 - k + m] \gamma^{\wedge} [P/2 - k + m] \gamma^{\alpha} + \gamma_{\alpha} [P/2 - k + m] \gamma^{\wedge} \frac{\partial}{\partial q_{\nu}} [(P-q)/2 - k + m] \gamma^{\alpha}$

with $\frac{\partial}{\partial q_{\nu}} [(P+q)/2 - k + m] = \frac{\gamma^{\wedge}}{2} \frac{\partial q_{\nu}}{\partial q_{\nu}} \gamma^{\wedge} = \frac{\gamma^{\nu}}{2}$

$$\Rightarrow \frac{\partial N^{\wedge}}{\partial q_{\nu}} \Big|_{q=0} = \gamma_{\alpha} \frac{\gamma^{\nu}}{2} \gamma^{\wedge} [P/2 - k + m] \gamma^{\alpha} - \gamma_{\alpha} [P/2 - k + m] \gamma^{\wedge} \frac{\gamma^{\nu}}{2} \gamma^{\alpha}$$

The derivative of the propagator is

$$\frac{\partial}{\partial q_\nu} \left(\frac{1}{[(P+q)/2 - k]^2 - k^2} [(P-q)/2 - k]^2 - k^2} \right)_{q=0}$$

$$= \frac{1}{[(P/2 - k)^2 - k^2]} \frac{\partial}{\partial q_\nu} \left(\frac{1}{[(P+q)/2 - k]^2 - k^2} \right)_{q=0} + \frac{1}{[(P/2 - k)^2 - k^2]} \frac{\partial}{\partial q_\nu} \left(\frac{1}{[(P-q)/2 - k]^2 - k^2} \right)_{q=0}$$

Now,

$$\begin{aligned} \frac{\partial}{\partial q_\nu} \left(\frac{1}{[(P \pm q)/2 - k]^2 - k^2} \right)_{q=0} &= \frac{1}{[(P/2 - k)^2 - k^2]^2} \cdot \frac{\partial}{\partial q_\nu} ((P \pm q)/2 - k)^2 \\ &= \frac{1}{[(P/2 - k)^2 - k^2]^2} \cdot 2(P/2 - k) \cdot \frac{1}{2} (\pm \delta_\nu^\nu) \\ &= \frac{\pm (P/2 - k)^\nu}{[(P/2 - k)^2 - k^2]^2} \end{aligned}$$

↑
Sign means that these terms
will sum to zero

$$\Rightarrow \frac{\partial}{\partial q_\nu} \left(\frac{1}{[(P+q)/2 - k]^2 - k^2} [(P-q)/2 - k]^2 - k^2} \right)_{q=0} = 0$$

So, the derivative of the vertex function is

$$\frac{\partial \Lambda^{\mu}}{\partial g_{\nu}} \Big|_{g=0} = -ie^2 \int \frac{d^4 k}{(2\pi)^4} \frac{1}{k^2 [(P/2 - k)^2 - m^2]^2} \frac{\partial N^{\mu}}{\partial g_{\nu}} \Big|_{g=0}$$

with,

$$\frac{\partial N^{\mu}}{\partial g_{\nu}} \Big|_{g=0} = \gamma_{\alpha} \gamma^{\nu} \gamma^{\mu} [P/2 - k + m] \gamma^{\alpha} - \gamma_{\alpha} [P/2 - k + m] \gamma^{\mu} \gamma^{\nu} \gamma^{\alpha}$$

In the projection formula, we need to take traces.

For the first term, we need

$$\tau_1 = \text{tr} \left[\left(\underbrace{m^2 \gamma_{\mu}}_{\text{red}} - \underbrace{P_{\mu} P}_{\text{blue}} - \underbrace{\frac{3}{2} m P_{\mu}}_{\text{green}} \right) N^{\mu}(P, 0, k) \right]$$

with

$$N^{\mu}(P, 0, k) = \gamma_{\alpha} [P/2 - k + m] \gamma^{\mu} [P/2 - k + m] \gamma^{\alpha}$$

So, the following compounds are

$$\gamma_{\alpha} \gamma_{\mu} \gamma^{\alpha} = -2\gamma_{\mu}$$

$$\begin{aligned} & \bullet \text{tr} \left[\gamma_{\mu} \gamma_{\alpha} [P/2 - k + m] \gamma^{\mu} [P/2 - k + m] \gamma^{\alpha} \right] \\ &= -2 \text{tr} \left[\gamma_{\mu} [P/2 - k + m] \gamma^{\mu} [P/2 - k + m] \right] \\ &= -2 \text{tr} \left[\gamma_{\mu} (P/2 - k) \gamma^{\mu} (P/2 - k) \right] - 2m^2 \text{tr} \left[\gamma_{\mu} \gamma^{\mu} \right] \\ &= +4 \text{tr} \left[(P/2 - k)(P/2 - k) \right] - 8m^2 \text{tr} \left[\mathbb{1} \right] \\ &= 16 (P/2 - k)^2 - 32m^2 \\ &= 16 \left(\frac{P^2}{4} + k^2 - P \cdot k \right) - 32m^2 \\ &= 16 (k^2 - P \cdot k - m^2) \end{aligned}$$

$$\begin{aligned}
& \bullet \operatorname{tr} \left[\mathcal{P} \gamma_\alpha [\mathcal{P}_{12} - k + m] \mathcal{P} [\mathcal{P}_{12} - k + m] \gamma^\alpha \right] \\
&= -2 \operatorname{tr} \left[\mathcal{P} [\mathcal{P}_{12} - k + m] \mathcal{P} [\mathcal{P}_{12} - k + m] \right] \\
&= -2 \operatorname{tr} \left[\mathcal{P} (\mathcal{P}_{12} - k) \mathcal{P} (\mathcal{P}_{12} - k) \right] - 2m^2 \operatorname{tr} [\mathcal{P} \mathcal{P}] \\
&= -8 \left[2 \mathcal{P} \cdot \left(\frac{\mathcal{P}}{2} - k \right) \mathcal{P} \cdot \left(\frac{\mathcal{P}}{2} - k \right) - \mathcal{P}^2 \left(\frac{\mathcal{P}}{2} - k \right)^2 \right] - 8m^2 \mathcal{P}^2 \\
&= -8 \left[2 \left(\frac{\mathcal{P}^2}{2} - k \cdot \mathcal{P} \right)^2 - 4m^2 \left(\frac{\mathcal{P}^2}{4} - k \cdot \mathcal{P} + k^2 \right) \right] - 32m^4 \\
&= -16 \left[(2m^2 - k \cdot \mathcal{P})^2 - 2m^2 (m^2 - k \cdot \mathcal{P} + k^2) \right] - 32m^4 \\
&= -16 \left[4m^4 + (k \cdot \mathcal{P})^2 - 4m^2 k \cdot \mathcal{P} - 2m^4 + 2m^2 k \cdot \mathcal{P} - 2m^2 k^2 + 2m^4 \right] \\
&= -16 \left[4m^4 + (k \cdot \mathcal{P})^2 - 2m^2 k \cdot \mathcal{P} - 2m^2 k^2 \right]
\end{aligned}$$

$$\begin{aligned}
& \bullet \operatorname{tr} \left[\gamma_\alpha [\mathcal{P}_{12} - k + m] \mathcal{P} [\mathcal{P}_{12} - k + m] \gamma^\alpha \right] \\
&= 4 \operatorname{tr} \left[(\mathcal{P}_{12} - k + m) \mathcal{P} (\mathcal{P}_{12} - k + m) \right] \\
&= 4m \operatorname{tr} \left[(\mathcal{P}_{12} - k) \mathcal{P} \right] + 4m \operatorname{tr} \left[\mathcal{P} (\mathcal{P}_{12} - k) \right] \\
&= 32m (\mathcal{P}_{12}^2 - k \cdot \mathcal{P}) \\
&= 32m (2m^2 - k \cdot \mathcal{P})
\end{aligned}$$

So, the first trace is

$$\begin{aligned}\tau_1 &= m^2 \cdot 16(k^2 - P \cdot k - m^2) + 16(4m^4 - 2m^2 P \cdot k - 2m^2 k^2 + (P \cdot k)^2) \\ &\quad - \frac{3}{2} m \cdot 32 m (2m^2 - k \cdot P) \\ &= -16m^2 k^2 + 16(P \cdot k)^2 - 48m^4 \\ &= -16[3m^4 + m^2 k^2 - (k \cdot P)^2]\end{aligned}$$

The second trace needed is

$$\tau_2 = \text{tr} \left[\left(\frac{P}{2} + m \right) [\gamma_r, \gamma_\nu] \left(\frac{P}{2} + m \right) \frac{\partial N^r}{\partial g_\nu} \Big|_{g=0} \right]$$

with

$$\begin{aligned}\frac{\partial N^r}{\partial g_\nu} \Big|_{g=0} &= \frac{1}{2} \left\{ \gamma_\alpha \gamma^\nu \gamma^\mu (P/2 - k + m) \gamma^\alpha \right. \\ &\quad \left. - \gamma_\alpha (P/2 - k + m) \gamma^\mu \gamma^\nu \gamma^\alpha \right\} \\ &= \frac{1}{2} \left\{ -2 \left(\frac{P}{2} - k \right) \gamma^\mu \gamma^\nu + 4m g^{\mu\nu} \right. \\ &\quad \left. + 2 \gamma^\nu \gamma^\mu \left(\frac{P}{2} - k \right) + 4m g^{\mu\nu} \right\}\end{aligned}$$

Note, $[\gamma_r, \gamma_\nu] g^{\mu\nu} = [\gamma_r, \gamma^\mu] = 0$

$$\begin{aligned}\text{So, } \tau_2 &= -\text{tr} \left[\left(\frac{P}{2} + m \right) \underline{[\gamma_r, \gamma_\nu]} \left(\frac{P}{2} + m \right) \left(\frac{P}{2} - k \right) \gamma^\mu \gamma^\nu \right] \\ &\quad + \text{tr} \left[\left(\frac{P}{2} + m \right) \underline{[\gamma_r, \gamma_\nu]} \left(\frac{P}{2} + m \right) \gamma^\nu \gamma^\mu \left(\frac{P}{2} - k \right) \right]\end{aligned}$$

$$\begin{aligned}
& \bullet \operatorname{tr}\left[\left(\frac{\not{P}}{2} + m\right)(\gamma_r \gamma_\nu - \gamma_\nu \gamma_r)\left(\frac{\not{P}}{2} + m\right)\left(\frac{\not{P}}{2} - k\right)\gamma^r \gamma^\nu\right] \\
&= m \operatorname{tr}\left[\frac{\not{P}}{2}(\gamma_r \gamma_\nu - \gamma_\nu \gamma_r)\left(\frac{\not{P}}{2} - k\right)\gamma^r \gamma^\nu\right] \\
&\quad + m \operatorname{tr}\left[(\gamma_r \gamma_\nu - \gamma_\nu \gamma_r)\frac{\not{P}}{2}\left(\frac{\not{P}}{2} - k\right)\gamma^r \gamma^\nu\right] \\
&= m \operatorname{tr}\left[\frac{\not{P}}{2}\gamma_r \gamma_\nu\left(\frac{\not{P}}{2} - k\right)\gamma^r \gamma^\nu\right] - m \operatorname{tr}\left[\frac{\not{P}}{2}\gamma_\nu \gamma_r\left(\frac{\not{P}}{2} - k\right)\gamma^r \gamma^\nu\right] \\
&\quad + m \operatorname{tr}\left[\gamma_r \gamma_\nu\frac{\not{P}}{2}\left(\frac{\not{P}}{2} - k\right)\gamma^r \gamma^\nu\right] - m \operatorname{tr}\left[\gamma_\nu \gamma_r\frac{\not{P}}{2}\left(\frac{\not{P}}{2} - k\right)\gamma^r \gamma^\nu\right] \\
&= 4m \operatorname{tr}\left[\frac{\not{P}}{2}\left(\frac{\not{P}}{2} - k\right)\right] + 2m \operatorname{tr}\left[\frac{\not{P}}{2}\gamma_\nu\left(\frac{\not{P}}{2} - k\right)\gamma^\nu\right] \\
&\quad - 8m \operatorname{tr}\left[\frac{\not{P}}{2}\left(\frac{\not{P}}{2} - k\right)\right] - 16m \operatorname{tr}\left[\frac{\not{P}}{2}\left(\frac{\not{P}}{2} - k\right)\right] \\
&= -20m \operatorname{tr}\left[\frac{\not{P}}{2}\left(\frac{\not{P}}{2} - k\right)\right] - 4m \operatorname{tr}\left[\frac{\not{P}}{2}\left(\frac{\not{P}}{2} - k\right)\right] \\
&= -24 \cdot 4m \frac{\not{P}}{2} \cdot \left(\frac{\not{P}}{2} - k\right) \\
&= -96m \left(\frac{\not{P}^2}{4} - k \cdot \not{P}\right) \\
&= -48m(2m^2 - k \cdot \not{P})
\end{aligned}$$

$$\begin{aligned}
& \bullet \operatorname{tr} \left[\left(\frac{\not{P} + \not{h}}{2} \right) (\gamma_r \gamma_\nu - \gamma_\nu \gamma_r) \left(\frac{\not{P} + \not{h}}{2} \right) \gamma^\nu \gamma^r \left(\frac{\not{P} - \not{h}}{2} \right) \right] \\
&= m \operatorname{tr} \left[\frac{\not{P}}{2} (\gamma_r \gamma_\nu - \gamma_\nu \gamma_r) \gamma^\nu \gamma^r \left(\frac{\not{P} - \not{h}}{2} \right) \right] \\
&\quad + m \operatorname{tr} \left[(\gamma_r \gamma_\nu - \gamma_\nu \gamma_r) \frac{\not{P}}{2} \gamma^\nu \gamma^r \left(\frac{\not{P} - \not{h}}{2} \right) \right] \\
&= m \operatorname{tr} \left[\frac{\not{P}}{2} \gamma_r \gamma_\nu \gamma^\nu \gamma^r \left(\frac{\not{P} - \not{h}}{2} \right) \right] - m \operatorname{tr} \left[\frac{\not{P}}{2} \gamma_\nu \gamma_r \gamma^\nu \gamma^r \left(\frac{\not{P} - \not{h}}{2} \right) \right] \\
&\quad + m \operatorname{tr} \left[\gamma_r \gamma_\nu \frac{\not{P}}{2} \gamma^\nu \gamma^r \left(\frac{\not{P} - \not{h}}{2} \right) \right] - m \operatorname{tr} \left[\gamma_\nu \gamma_r \frac{\not{P}}{2} \gamma^\nu \gamma^r \left(\frac{\not{P} - \not{h}}{2} \right) \right] \\
&= 16 m \operatorname{tr} \left[\frac{\not{P}}{2} \left(\frac{\not{P} - \not{h}}{2} \right) \right] + 8 m \operatorname{tr} \left[\frac{\not{P}}{2} \left(\frac{\not{P} - \not{h}}{2} \right) \right] \\
&\quad + 4 m \operatorname{tr} \left[\frac{\not{P}}{2} \left(\frac{\not{P} - \not{h}}{2} \right) \right] - 4 m \operatorname{tr} \left[\frac{\not{P}}{2} \left(\frac{\not{P} - \not{h}}{2} \right) \right] \\
&= 24 \cdot 4 m \frac{\not{P}}{2} \cdot \left(\frac{\not{P} - \not{h}}{2} \right) \\
&= 96 m \left(\frac{\not{P}^2}{4} - \not{P} \cdot \not{h} \right) \\
&= 48 m (2m^2 - 4 \cdot P)
\end{aligned}$$

$$\text{So, } \tau_2 = 96 m (2m^2 - 4 \cdot P)$$

Therefore, we find for the anomaly

$$a_e^{(1)} = \frac{1}{12m^2} \text{tr} \left[\left(m^2 \gamma_\mu - \not{P} \not{P} - \frac{3}{2} m \not{P}_\mu \right) V_1^\mu \right. \\ \left. + \frac{m}{4} \left(\not{P} + m \right) [\gamma_\mu, \gamma_\nu] \left(\not{P} + m \right) \delta V_1^{\mu\nu} \right]$$

$$= -\frac{i e^2}{12m^2} \int \frac{d^4 k}{(2\pi)^4} \frac{1}{k^2 [(P/2 - k)^2 - m^2]^2} \cdot \left[\tau_1 + \frac{m}{4} \tau_2 \right]$$

Now,

$$\tau_1 + \frac{m}{4} \tau_2 = -16 \left[3m^4 + m^2 k^2 - (k \cdot P)^2 \right] \\ + 24 m^2 (2m^2 - k \cdot P)$$

$$= -16 m^2 k^2 + 16 (k \cdot P)^2 - 24 m^2 k \cdot P$$

Notice that the integrand is finite as $k \rightarrow 0$, since the numerator is proportional to k . Naively, we see that the integral could be UV divergent, however we will see that this integral is finite.

So,

$$a_e^{(1)} = -i e^2 \int \frac{d^4 k}{(2\pi)^4} \frac{1}{k^2 [(P/2 - k)^2 - m^2]^2} \left(-\frac{4}{3} k^2 + \frac{4}{3m^2} (k \cdot P)^2 - 2k \cdot P \right)$$

We now introduce Feynman parameters

$$\begin{aligned} \frac{1}{A^2 B} &= 2 \int_0^1 dx_1 \int_0^1 dx_2 \delta(1-x_1-x_2) \frac{x_1}{[x_1 A + x_2 B]^3} \\ &= 2 \int_0^1 dx \frac{x}{[x A + (1-x) B]^3} \\ &= 2 \int_0^1 dx \frac{x}{[x(A-B) + B]^3} \end{aligned}$$

so, with

$$A = (P/2 - k)^2 - m^2 = \frac{P^2}{4} + k^2 - k \cdot P - m^2$$

and

$$B = k^2$$

$$= \frac{4}{4} k^2 - k \cdot P$$

so,

$$\frac{1}{k^2 [(P/2 - k)^2 - m^2]^2} = 2 \int_0^1 dx \frac{x}{[k^2 - k \cdot P x]^3}$$

so,

$$a_e^{(1)} = -2ie^2 \int_0^1 dx x \int \frac{d^4 k}{(2\pi)^4} \frac{1}{(k^2 - k \cdot P x)^3} \left(-\frac{4}{3} k^2 + \frac{4}{3k^2} (k \cdot P)^2 - 2k \cdot P \right)$$

To perform the momentum integral, let us perform this in d -dimensions to see the cancellation of the UV divergence.

Note the following integrals (No proof)

$$I(n) = \int d^d k \frac{1}{(k^2 + 2\hat{p} \cdot k - M^2)^n} = i\pi^{d/2} \frac{(-1)^n}{\Gamma(n)} \frac{\Gamma(n-d/2)}{(\hat{p}^2 + M^2)^{n-d/2}}$$

$$I_\mu(n) = \int d^d k \frac{k_\mu}{(k^2 + 2\hat{p} \cdot k - M^2)^n} = -\hat{p}_\mu I(n)$$

$$I_{\mu\nu}(n) = \int d^d k \frac{k_\mu k_\nu}{(k^2 + 2\hat{p} \cdot k - M^2)^n} = \left[\hat{p}_\mu \hat{p}_\nu - \frac{1}{2} g_{\mu\nu} \left(\frac{\hat{p}^2 + M^2}{n-1-d/2} \right) \right] I(n)$$

For us, $\hat{p} = -\frac{P_x}{2}$, $M^2 = 0$, $n = 3$, and $\text{with } d = 4 - 2\epsilon$

$$\text{so } I(3) = -\frac{i\pi^2}{2\hat{p}^2} = -\frac{2i\pi^2}{P_x^2} = -\frac{i\pi^2}{2m^2 x^2}$$

↑
take $\epsilon \rightarrow 0$
at end

and

$$I_\mu(3) = +\frac{x}{2} P_\mu I(3)$$

$$\begin{aligned} I_{\mu\nu}(3) &= \left[\frac{x^2}{4} P_\mu P_\nu - \frac{1}{2} g_{\mu\nu} \frac{(P_x^2/4)}{2-d/2} \right] I(3) \\ &= \frac{x^2}{4} \left[P_\mu P_\nu - \frac{2m^2}{\epsilon} g_{\mu\nu} \right] I(3) \end{aligned}$$

Therefore,

$$\begin{aligned} \int d^4 k \frac{1}{(k^2 - k \cdot P_x)^3} \left(-\frac{4}{3} k^2 + \frac{4}{3m^2} (k \cdot P)^2 - 2k \cdot P \right) \\ = -\frac{4}{3} g^{\mu\nu} I_{\mu\nu}(3) + \frac{4}{3m^2} P^\mu P^\nu I_\mu(3) - 2P^\mu I_\mu(3) \end{aligned}$$

$$\begin{aligned}
& \int d^4k \frac{1}{(k^2 - k \cdot P_x)^3} \left(-\frac{4}{3} k^2 + \frac{4}{3h^2} (k \cdot P)^2 - 2k \cdot P \right) \\
&= -\frac{4}{3} g^{\mu\nu} \underline{I}_{\mu\nu}(3) + \frac{4}{3h^2} P^\mu P^\nu \underline{I}_{\mu\nu}(1) - 2P^\mu \underline{I}_\mu(3) \\
&= -\frac{4}{3} g^{\mu\nu} \frac{x^2}{4} \left[\underline{P}_\mu \underline{P}_\nu - \frac{2m^2}{\epsilon} g_{\mu\nu} \right] \underline{I}(3) \\
&\quad + \frac{4}{3h^2} P^\mu P^\nu \frac{x^2}{4} \left[\underline{P}_\mu \underline{P}_\nu - \frac{2m^2}{\epsilon} g_{\mu\nu} \right] \underline{I}(3) \\
&\quad - 2P^\mu \frac{x}{2} \underline{P}_\mu \underline{I}(3) \\
&= \left\{ -\frac{4}{3} \frac{x^2}{4} \left[P^2 - \frac{2m^2}{\epsilon} \cdot 4 \right] + \frac{4}{3h^2} \frac{x^2}{4} \left[P^2 P^2 - \frac{2m^2}{\epsilon} P^2 \right] - x P^2 \right\} \underline{I}(3) \\
&= \left\{ \frac{x^2}{3} \left[-P^2 + \frac{2m^2}{\epsilon} \cdot 4 \right] + \frac{P^2 P^2}{h^2} - \frac{2m^2}{\epsilon} \frac{P^2}{h^2} \right\} \underline{I}(3) \\
&= \left\{ \frac{x^2}{3} \left[-4m^2 + 16m^2 \right] - x \cdot 4m^2 \right\} \underline{I}(3) \quad \text{Naive UV divergence cancels.} \\
&= 4m^2 x (x-1) \underline{I}(3) \\
&= -2 \frac{(x-1)}{x} i\pi^2
\end{aligned}$$

Therefore, the simplified expression for $a_e^{(1)}$ is

$$\begin{aligned}
a_e^{(1)} &= -2ie^2 \int_0^1 dx \, x \left\{ \frac{2}{x} (1-x) i\pi^2 \right\} \\
&= + \frac{(2\pi)^2}{(2\pi)^4} e^2 \int_0^1 dx (1-x) = \frac{e^2}{4\pi^2} \left[x - \frac{x^2}{2} \right]_0^1 = \frac{1}{2} \left(\frac{e^2}{4\pi^2} \right)
\end{aligned}$$

So,

$$a_e^{(1)} = \frac{1}{2} \left(\frac{e^2}{4\pi^2} \right)$$

Recall the fine-structure constant, $\alpha \equiv \frac{e^2}{4\pi^2}$

$$\Rightarrow \boxed{a_e^{(1)} = \frac{\alpha}{2\pi}} \quad \text{Schwinger's triumph!}$$

here, $a_e^{(1)} = 1.16 \dots \times 10^{-3}$

The measured value is $(a_e)^{\text{exp}} = 1.15965218073(28) \times 10^{-3}$

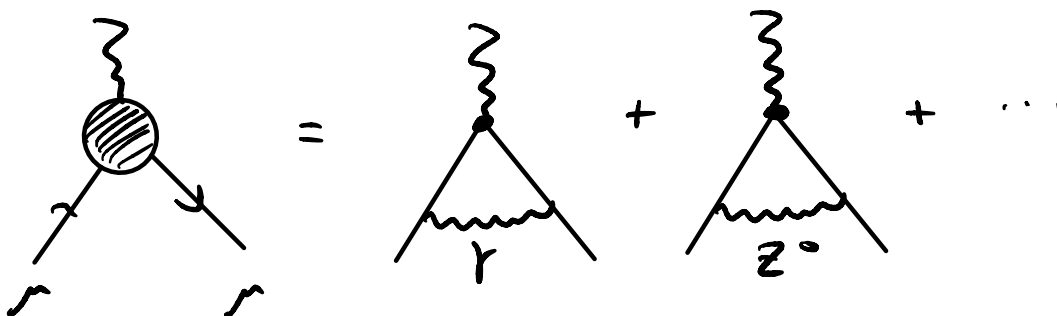
While the current theoretical prediction ($O(\alpha^{10})!$) is

$$(a_e)^{\text{SM}} = 1.15965218188(78) \times 10^{-3}$$

A remarkable prediction!

For muons, the leading order is identical, $a_\mu^{(1)} = \frac{\alpha}{2\pi}$, but since $\frac{m_\mu}{m_e} \sim 200$, the value of $\frac{g-2}{2}$ is susceptible

to searches for new physics.



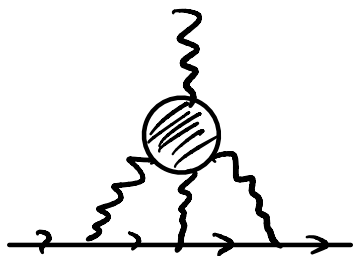
EW coupling.

For the Z⁰-boson, $\alpha_\mu \sim \left(\frac{m_\mu}{m_Z}\right)^2 \cdot \frac{g^2}{16\pi} \cdot \log\left(\frac{m_Z^2}{m_\mu^2}\right)$

There is an intense search for new physics by looking at the muon anomaly. The current (as of Aug. 2023) experimental measurement is

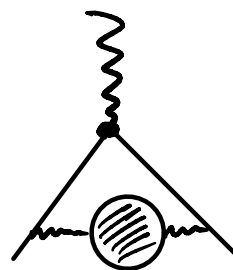
$$(\alpha_\mu)^{exp} = 116592059(22) \times 10^{-11}$$

The SM theory is a little more muddled. There are discrepancies between lattice methods and data-driven analyses. There is an apparent 4.2σ discrepancy with experiment. BQ, two hadronic processes are the leading candidates



Hadronic light-by-light

and



Hadronic vacuum polarization