

PHYS 772 – The Standard Model of Particle Physics

Problem Set 4 – Solution

Due: Tuesday, March 04 at 4:00 pm

Term: Spring 2025

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1. Show that the Lie algebra structure constants c_{jkl} , defined by the Lie bracket $[X^j, X^k] = c_{jkl}X^l$, satisfy the relation $c_{jkm}c_{mln} + c_{klm}c_{mjn} + c_{ljm}c_{mkn} = 0$.

Solution: Here we use the Jacobi identity for the elements of the Lie algebra,

$$\sum_{(j,k,l)} [[X^j, X^k], X^l] = 0,$$

where (j, k, l) indicates the cyclic sum. Performing the sum, and using $[X^j, X^k] = c_{jkl}X^l$, we find

$$\sum_{(j,k,l)} [[X^j, X^k], X^l] = [[X^j, X^k], X^l] + [[X^k, X^l], X^j] + [[X^l, X^j], X^k],$$

$$= c_{jkm} [X^m, X^l] + c_{klm} [X^m, X^j] + c_{ljm} [X^m, X^k],$$

$$= c_{jkm} c_{mln} X^n + c_{klm} c_{mjn} X^n + c_{ljm} c_{mkn} X^n,$$

$$= (c_{jkm} c_{mln} + c_{klm} c_{mjn} + c_{ljm} c_{mkn}) X^n,$$

$$= 0.$$

Since this must be zero for any X^n , we must have $c_{jkm}c_{mln} + c_{klm}c_{mjn} + c_{ljm}c_{mkn} = 0$, as desired.

2. Consider a general Lie algebra $[X^j, X^k] = c_{jkl}X^l$, where $c_{jkl} = -c_{kjl}$. From the structure constants, we may form matrices M^j with matrix elements $(M^j)_{lk} = c_{jkl}$. Note the order of the indices. Show that these matrices furnish a representation of the algebra, i.e., show that $[M^j, M^k] = c_{jkl}M^l$. This representation is called the *adjoint representation*. **Hint:** The Jacobi identity may be helpful.

Solution: Looking at the matrix elements of the commutator of $(M^j)_{lk} = c_{jkl}$, $([M^j, M^k])_{ln} = (M^j)_{lm} (M^k)_{mn} - (M^k)_{lm} (M^j)_{mn}$, $= c_{jml}c_{knm} - c_{kml}c_{jnm}$, $= c_{ljm}c_{mkn} + c_{klm}c_{mjn}$,

where we used the antisymmetry of the structure constants. From the Jacobi identity of the structure constants, $c_{jkm}c_{mln} + c_{klm}c_{mjn} + c_{ljm}c_{mkn} = 0$, we have $c_{ljm}c_{mkn} + c_{klm}c_{mjn} =$

 $-c_{jkm}c_{mln}$. So, the commutator is

$$([M^{j}, M^{k}])_{ln} = c_{ljm}c_{mkn} + c_{klm}c_{mjn},$$
$$= -c_{jkm}c_{mln},$$
$$= c_{jkm}c_{mnl},$$
$$= c_{jkm}(M^{m})_{ln}.$$

Therefore $[M^j, M^k] = c_{jkl}M^l$, and we conclude that $(M^j)_{lk} = c_{jkl}$ is a valid representation of the Lie algebra.

3. Suppose X^j is a generator for the Lie algebra $[X^j, X^k] = c_{jkl}X^l$. Show that $X^2 = \sum_j X^j X^j$ commutes with the group generators, and therefore we may write $(X^2)_{ab} = C_2(r) \delta_{ab}$ where $C_2(r)$ is a constant called the *quadratic Casimir* of the representation r.

Solution: We want to show that $[X^2, X^k] = 0$ where $X^2 = \sum_j X^j X^j$ and $[X^j, X^k] = c_{jkl} X^l$. So, taking the commutator

$$\begin{split} [X^2, X^k] &= \sum_j [X^j X^j, X^k] \,, \\ &= \sum_j X^j [X^j, X^k] + \sum_j [X^j, X^k] X^j \end{split}$$

where we used [AB, C] = A[B, C] + [A, C]B. Now, we use $[X^j, X^k] = c_{jkl}X^l$, noting l is being summed over implicitly. So, we find

$$\begin{split} X^{2}, X^{k}] &= \sum_{j} X^{j} [X^{j}, X^{k}] + \sum_{j} [X^{j}, X^{k}] X^{j} ,\\ &= \sum_{j,l} X^{j} (c_{jkl} X^{l}) + \sum_{j,l} (c_{jkl} X^{l}) X^{j} ,\\ &= \sum_{j,l} c_{jkl} X^{j} X^{l} + \sum_{j,l} c_{jkl} X^{l} X^{j} ,\\ &= \sum_{j,l} c_{jkl} X^{j} X^{l} + \sum_{j,l} c_{lkj} X^{j} X^{l} ,\\ &= \sum_{j,l} c_{jkl} X^{j} X^{l} - \sum_{j,l} c_{jkl} X^{j} X^{l} = 0 ,\end{split}$$

where in the fourth line we interchanged the summed indices, and in going to the last line we noted that $c_{lkj} = -c_{jkl}$. Therefore, X^2 commutes with all the generators X^j . Therefore, we can write $(X^2)_{ab} = C_2(r)\delta_{ab}$, where $C_2(r)$ is some constant which depends on the representation r, and a, b span the dimension of the representation, $a, b = 1, \ldots, r$.

4. Let X^j be a generator for a generic $\mathfrak{su}(N)$ Lie algebra, $[X^j, X^k] = c_{jkl}X^l$, and $U(\alpha^j)$ is an element

of the corresponding Lie group SU(N), with $U(\alpha^j) = \exp(\alpha^j X_j)$ with $\alpha^j \in \mathbb{R}$. Show that X^j are traceless, antihermitian $N \times N$ matrices.

Solution: Since $U(\alpha^j) \in SU(N)$, then we require

 $U(\alpha^j)^{\dagger}U(\alpha^j) = U(\alpha^j)U(\alpha^j)^{\dagger} = I_N,$

where I_N is the $N \times N$ identity. Furthermore, $\det(U(\alpha^j)) = 1$. From the properties of matrix exponentials, $\exp(\alpha^j X_j)^{\dagger} = \exp(\alpha^j X_j^{\dagger})$. Let us Taylor expand the product $U(\alpha^j)^{\dagger}U(\alpha^j)$ about $\alpha^j = 0$,

$$I_N = U(\alpha^j)^{\dagger} U(\alpha^j) = \left(I_N + \alpha^j X_j^{\dagger} + \mathcal{O}(\alpha^2) \right) \left(I_N + \alpha^j X_j + \mathcal{O}(\alpha^2) \right) ,$$

$$= I_N + \alpha^j X_j + \alpha^j X_j^{\dagger} + \mathcal{O}(\alpha^2) ,$$

$$= I_N + \alpha^j (X_j + X_j^{\dagger}) + \mathcal{O}(\alpha^2) ,$$

Since this must hold order-by-order in α , we have $X_j + X_j^{\dagger} = 0$, or $X_j = -X_j^{\dagger}$, proving that the generators are antihermitian. Next, recall for matrix exponentials $\det(\exp^A) = \exp(\operatorname{tr}(A))$ where A is an $N \times N$ matrix. Since $\det(U(\alpha^j)) = 1$, we have the following

$$1 = \det(U(\alpha^j)) = \det\left(\exp(\alpha^j X_j)\right),$$
$$= \exp\left(\operatorname{tr}(\alpha^j X_j)\right),$$
$$= \exp\left(\alpha^j \operatorname{tr}(X_j)\right).$$

Since this must hold for any α^j , we conclude that $tr(X_j) = 0$.

5. Consider the set of all complex 2×2 matrices M with det(M) = i. Does this set form a group under the usual matrix multiplication? Explain your reasoning.

Solution: Let G be the set of all 2×2 matrices with $\det(M) = i$. Let us assume that $M \in G$, some group where $\det(M) = i$. If M_1 and M_2 are elements of the group, then the product $M_1 \cdot M_2$ should close under the group multiplication, that is $M_3 = M_1 \cdot M_2 \in G$. Since M_3 is in G, then $\det(M_3) = i$. But, consider $\det(M_3) = \det(M_1 \cdot M_2) = \det(M_1) \det(M_2)$, from the properties of determinants. So, $\det(M_3) = \det(M_1) \det(M_2) = (i)(i) = -1 \neq i$, which contradicts our assumption. Therefore, the product of two group elements does not close under group multiplication, and thus G does not form a group.

Alternatively, assume the existence of an inverse matrix $M^{-1} \in G$. The determinant of an inverse matrix is $\det(M^{-1}) = 1/\det(M) = 1/i = -i \neq i$. Therefore, we conclude that such an inverse matrix does not exist, and therefore G is not a group.

6. Consider $X_j = -\frac{1}{2}i\sigma_j$ as a bases element of the $\mathfrak{su}(2)$ algebra, $[X_j, X_k] = \epsilon_{jkl}X_l$, where σ_j are the

Pauli matrices,

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \qquad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \qquad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

Verify the following:

(a) $[\sigma_j, \sigma_k] \equiv \sigma_j \sigma_k - \sigma_k \sigma_j = 2i\epsilon_{jkl}\sigma_l$.

Solution: We compute the following products,

$$\sigma_{1}\sigma_{2} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} = i\sigma_{3},$$

$$\sigma_{2}\sigma_{1} = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} -i & 0 \\ 0 & i \end{pmatrix} = -i\sigma_{3},$$

$$\sigma_{2}\sigma_{3} = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} = i\sigma_{1},$$

$$\sigma_{3}\sigma_{2} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \begin{pmatrix} 0 & -i \\ -i & 0 \end{pmatrix} = -i\sigma_{1}$$

$$\sigma_{3}\sigma_{1} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = i\sigma_{2},$$

$$\sigma_{1}\sigma_{3} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} = -i\sigma_{2},$$

as well as the squares of the Pauli matrices

$$\sigma_1^2 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = I_2,$$

$$\sigma_2^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = I_2,$$

$$\sigma_3^2 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = I_2.$$

Therefore, we have $\sigma_j^2 = I_2$ and $\sigma_j \sigma_k = -\sigma_k \sigma_j$ for $j \neq k$. So, the commutator $[\sigma_j, \sigma_j] = 0$, while $[\sigma_1, \sigma_2] = +2i\sigma_3$, $[\sigma_2, \sigma_3] = +2i\sigma_1$, and $[\sigma_3, \sigma_1] = +2i\sigma_2$. The commutators are completely antisymmetric, thus we can write it in terms of the Levi-Civita ϵ_{jkl} tensor, $[\sigma_j, \sigma_k] = 2i\epsilon_{jkl}\sigma_l$.

(b) $\{\sigma_j, \sigma_k\} \equiv \sigma_j \sigma_k + \sigma_k \sigma_j = 2\delta_{jk}I_2.$

Solution: From the results of part (a), we find that $\sigma_j \sigma_k + \sigma_k \sigma_j = 0$ for $j \neq k$. Therefore, $\{\sigma_k, \sigma_k\} = \sigma_j \sigma_k + \sigma_k \sigma_j = \delta_{jk} (2\sigma_j \sigma_j) = 2\delta_{jk} I_2$.

(c) $\sigma_j \sigma_k = \delta_{jk} I_2 + i \epsilon_{jkl} \sigma_l.$

Solution: Let us add the two results $[\sigma_j, \sigma_k] = 2i\epsilon_{jkl}\sigma_l$ and $\{\sigma_j, \sigma_k\} = 2\delta_{jk}I_2$,

$$[\sigma_j, \sigma_k] + \{\sigma_j, \sigma_k\} = 2\sigma_j \sigma_k = 2i\epsilon_{jkl}\sigma_l + 2\delta_{jk}I_2.$$

Therefore, we immediately find that $\sigma_j \sigma_k = \delta_{jk} I_2 + i \epsilon_{jkl} \sigma_l$.

(d) Show that a group element $U(\alpha^j) \in SU(2)$ can be written as

$$U(\alpha^{j}) = \exp\left(-\frac{1}{2}i\alpha^{j}\sigma_{j}\right) = I_{2}\cos\left(\frac{1}{2}\alpha\right) - i\frac{\alpha^{j}\sigma_{j}}{\alpha}\sin\left(\frac{1}{2}\alpha\right),$$

where $\alpha^2 = \sum_j (\alpha_j)^2$.

Solution: Let us Taylor expand about $\alpha^j = 0$, $\exp\left(-\frac{1}{2}i\alpha^j\sigma_j\right) = \sum_{n=0}^{\infty} \frac{1}{n!} \left(-\frac{1}{2}i\alpha^j\sigma_j\right)^n,$ $= \sum_{n=0}^{\infty} \frac{1}{(2n)!} \left(-\frac{1}{2}i\alpha^j\sigma_j\right)^{2n} + \sum_{n=0}^{\infty} \frac{1}{(2n+1)!} \left(-\frac{1}{2}i\alpha^j\sigma_j\right)^{2n+1},$

where we split the sum into even and odd terms. Now, $(-i\alpha^j\sigma_j/2)^{2n} = (-1)^n (\alpha^j\sigma_j/2)^{2n}$, while $(-i\alpha^j\sigma_j/2)^{2n+1} = -i(-1)^n (\alpha^j\sigma_j/2)^{2n+1}$. Now, we evaluate $(\alpha^j\sigma_j)^2$,

$$\begin{aligned} \alpha^{j}\sigma_{j})^{2} &= (\alpha^{j}\sigma_{j})(\alpha^{k}\sigma_{k}), \\ &= \alpha^{j}\alpha^{k}\left(\sigma_{j}\sigma_{k}\right), \\ &= \alpha^{j}\alpha^{k}\left(\delta_{jk}I_{2} + i\epsilon_{jkl}\sigma_{l}\right), \\ &= \alpha^{j}\alpha^{j}I_{2} = \alpha^{2}I_{2}. \end{aligned}$$

So, $(\alpha^j \sigma_j)^{2n} = (\alpha)^{2n} I_2$, and $(\alpha^j \sigma_j)^{2n+1} = (\alpha^j \sigma_j)^{2n} (\alpha^j \sigma_j) = (\alpha)^{2n} (\alpha^j \sigma_j)$. So, the exponential expansion is

$$\exp\left(-\frac{1}{2}i\alpha^{j}\sigma_{j}\right) = \sum_{n=0}^{\infty} \frac{(-1)^{n}}{(2n)!} \left(\frac{1}{2}\alpha^{j}\sigma_{j}\right)^{2n} - i\sum_{n=0}^{\infty} \frac{(-1)^{n}}{(2n+1)!} \left(\frac{1}{2}\alpha^{j}\sigma_{j}\right)^{2n+1},$$
$$= I_{2}\sum_{n=0}^{\infty} \frac{(-1)^{n}}{(2n)!} \left(\frac{\alpha}{2}\right)^{2n} - i\left(\frac{\alpha^{j}\sigma_{j}}{\alpha}\right) \sum_{n=0}^{\infty} \frac{(-1)^{n}}{(2n+1)!} \left(\frac{\alpha}{2}\right)^{2n+1},$$
$$= I_{2}\cos\left(\frac{\alpha}{2}\right) - i\frac{\alpha^{j}\sigma_{j}}{\alpha}\sin\left(\frac{\alpha}{2}\right).$$

7. Consider $X_j = L_j$ as a bases element of the $\mathfrak{so}(3)$ algebra, $[X_j, X_k] = \epsilon_{jkl} X_l$, where L_j are the matrices,

$$L_1 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix}, \qquad L_2 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ -1 & 0 & 0 \end{pmatrix}, \qquad L_3 = \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

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Verify the following:

(a)
$$[L_j, L_k] = \epsilon_{jkl} L_l$$
.

Solution: Notice that $(L^j)_{lk} = \epsilon_{jkl}$. So, the commutator is

$$([L^{j}, L^{k}])_{ln} = (L^{j})_{lm} (L^{k})_{mn} - (L^{k})_{lm} (L^{j})_{mn},$$

 $=\epsilon_{jml}\epsilon_{knm}-\epsilon_{kml}\epsilon_{jnm}\,.$

Recall the Jacobi identity for the structure constants, here the Levi-Civita, $\epsilon_{jkm}\epsilon_{mln} + \epsilon_{klm}\epsilon_{mjn} + \epsilon_{ljm}\epsilon_{mkn} = 0$. We can use the antisymmetry of ϵ_{jkl} to write

$$-\epsilon_{jkm}\epsilon_{mln} = \epsilon_{klm}\epsilon_{mjn} + \epsilon_{ljm}\epsilon_{mkn} ,$$

$$= \epsilon_{jml}\epsilon_{knm} - \epsilon_{kml}\epsilon_{jnm} ,$$

$$= (L^j)_{lm}(L^k)_{mn} - (L^k)_{lm}(L^j)_{mn} ,$$

where we identified the difference in Levi-Civita's as the commutator of $[L^j, L^k]$. So, we find

$$(L^{j})_{lm}(L^{k})_{mn} - (L^{k})_{lm}(L^{j})_{mn} = -\epsilon_{jkm}\epsilon_{mln},$$
$$= \epsilon_{jkm}\epsilon_{mnl},$$
$$= \epsilon_{jkm}(L^{m})_{ln}$$

We conclude that $[L^j, L^k] = \epsilon_{jkl} L^l$.

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(b) $\{L_j, L_k\} \neq N\delta_{jk}$ for any j, k, and N.

Solution: The anticommutator $\{L_j, L_k\} = L_j L_k + L_k L_j$. Since $(L^j)_{lk} = \epsilon_{jkl}$, we have $(L^j)_{lm}(L^k)_{mn} + (L^k)_{lm}(L^j)_{mn} = \epsilon_{jml}\epsilon_{knm} + \epsilon_{kml}\epsilon_{jnm}$, $= -\epsilon_{jlm}\epsilon_{knm} - \epsilon_{klm}\epsilon_{jnm}$,

where in the last line we used the antisymmetry properties of the permutation tensor. Now, we use the property $\epsilon_{jlm}\epsilon_{knm} = \delta_{jk}\delta_{ln} - \delta_{jn}\delta_{kl}$. So, we have

$$(L^{j})_{lm}(L^{k})_{mn} + (L^{k})_{lm}(L^{j})_{mn} = -\epsilon_{jlm}\epsilon_{knm} - \epsilon_{klm}\epsilon_{jnm},$$

$$= -(\delta_{jk}\delta_{ln} - \delta_{jn}\delta_{kl}) - (\delta_{kj}\delta_{ln} - \delta_{kn}\delta_{lj}),$$

$$= -2\delta_{jk}\delta_{ln} + \delta_{jn}\delta_{kl} + \delta_{kn}\delta_{lj}.$$

Thus, we see that $\{L_j, L_k\} \neq N\delta_{jk}$ for any j, k, and N.

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(c) Show that a group element $O(\alpha^j) \in SO(3)$ can be written as

$$O(\alpha^{j}) = \exp\left(\alpha^{j}L_{j}\right) = I_{3} + \frac{\alpha^{j}L_{j}}{\alpha}\sin\alpha + \left(\frac{\alpha^{j}L_{j}}{\alpha}\right)^{2}\left(1 - \cos\alpha\right),$$

where $\alpha^2 = \sum_j (\alpha_j)^2$.

Solution: Taylor expanding about $\alpha^j = 0$, we find

$$\exp\left(\alpha^{j}L_{j}\right) = \sum_{n=0}^{\infty} \frac{1}{n!} \left(\alpha^{j}L_{j}\right)^{n},$$
$$= \sum_{n=0}^{\infty} \frac{1}{(2n)!} \left(\alpha^{j}L_{j}\right)^{2n} + \sum_{n=0}^{\infty} \frac{1}{(2n+1)!} \left(\alpha^{j}L_{j}\right)^{2n+1},$$

where we split the sum into even and odd terms. Let us evaluate $(\alpha^j L_j)^2$,

$$[(\alpha^{j}L_{j})^{2}]_{ln} = [(\alpha^{j}L_{j})(\alpha^{k}L_{k})]_{ln},$$

$$= \alpha^{j}\alpha^{k} (L_{j})_{lm}(L_{k})_{mn},$$

$$= \alpha^{j}\alpha^{k} \epsilon_{jml}\epsilon_{knm},$$

$$= -\alpha^{j}\alpha^{k} \epsilon_{jlm}\epsilon_{knm},$$

$$= -\alpha^{j}\alpha^{k} (\delta_{jk}\delta_{ln} - \delta_{jn}\delta_{kl}),$$

$$= -(\alpha^{2}\delta_{ln} - \alpha_{n}\alpha_{l}),$$

where in the fourth line we used $\epsilon_{jml} = -\epsilon_{jlm}$, and in the fifth line we used the property $\epsilon_{jlm}\epsilon_{knm} = \delta_{jk}\delta_{ln} - \delta_{jn}\delta_{kl}$. Next, we evaluate $(\alpha^j L_j)^3$,

$$[(\alpha^{j}L_{j})^{3}]_{lp} = [(\alpha^{j}L_{j})(\alpha^{k}L_{k})(\alpha^{r}L_{r})]_{lp},$$

$$= \alpha^{j}\alpha^{k}\alpha^{r}(L_{j})_{lm}(L_{k})_{mn}(L_{r})_{np},$$

$$= -(\alpha^{2}\delta_{ln} - \alpha^{n}\alpha^{l})\alpha^{r}(L_{r})_{np},$$

$$= -\alpha^{2}\alpha^{r}(L_{r})_{lp} + \alpha^{n}\alpha^{l}\alpha^{r}(L_{r})_{np}$$

$$= -\alpha^{2}\alpha^{r}(L_{r})_{lp},$$

$$= -\alpha^{2}(\alpha^{j}L_{j})_{lp},$$

where we used that $\alpha^n \alpha^r (L_r)_{np} = \alpha^n \alpha^r \epsilon_{rpn} = 0$ since we have a symmetric sum over a completely antisymmetric object. So, we find that $(\alpha^j L_j)^3 = -\alpha^2 (\alpha^j L_j)$. Finally, we evaluate $(\alpha^j L_j)^4$ as

$$(\alpha^j L_j)^4 = (\alpha^j L_j)^3 (\alpha^j L_j),$$

= $-\alpha^2 (\alpha^k L_k) (\alpha^j L_j) = -\alpha^2 (\alpha^j L_j)^2.$

Therefore, the sum over even terms can be expressed in terms of $(\alpha^j L_j)^2$, while the odd terms can be written in terms of $(\alpha^j L_j)$. Specifically,

$$(\alpha^{j}L_{j})^{2n} = (-1)^{n+1}\alpha^{2n} \left(\frac{\alpha^{j}L_{j}}{\alpha}\right)^{2}, \quad n > 0$$
$$(\alpha^{j}L_{j})^{2n+1} = (-1)^{n}\alpha^{2n+1} \left(\frac{\alpha^{j}L_{j}}{\alpha}\right), \quad n \ge 0.$$

Substituting these expressions into the series expansion,

$$\begin{split} \exp\left(\alpha^{j}L_{j}\right) &= I_{3} + \sum_{n=1}^{\infty} \frac{1}{(2n)!} \left(\alpha^{j}L_{j}\right)^{2n} + \sum_{n=0}^{\infty} \frac{1}{(2n+1)!} \left(\alpha^{j}L_{j}\right)^{2n+1}, \\ &= I_{3} - \left(\frac{\alpha^{j}L_{j}}{\alpha}\right)^{2} \sum_{n=1}^{\infty} \frac{(-1)^{n}}{(2n)!} \alpha^{2n} + \left(\frac{\alpha^{j}L_{j}}{\alpha}\right) \sum_{n=0}^{\infty} \frac{(-1)^{n}}{(2n+1)!} \alpha^{2n+1}, \\ &= I_{3} + \left(\frac{\alpha^{j}L_{j}}{\alpha}\right)^{2} \left(1 - \sum_{n=0}^{\infty} \frac{(-1)^{n}}{(2n)!} \alpha^{2n}\right) + \left(\frac{\alpha^{j}L_{j}}{\alpha}\right) \sum_{n=0}^{\infty} \frac{(-1)^{n}}{(2n+1)!} \alpha^{2n+1}, \\ &= I_{3} + \left(\frac{\alpha^{j}L_{j}}{\alpha}\right)^{2} \left(1 - \cos\alpha\right) + \left(\frac{\alpha^{j}L_{j}}{\alpha}\right) \sin\alpha \,. \end{split}$$